

# When Is Water Not Water? Exploring Water Confined in Large Reverse Micelles Using a Highly Charged Inorganic Molecular Probe

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Abstract: The interior water pool of aerosol OT (AOT) reverse micelles tends toward bulk water properties as the micelle size increases. Thus, deviations from bulk water behavior in large reverse micelles are less expected than in small reverse micelles. Probing the interior water pool of AOT reverse micelles with a highly charged decavanadate (V<sub>10</sub>) oligomer using <sup>51</sup>V NMR spectroscopy shows distinct changes in solute environment. For example, when an acidic stock solution of protonated V<sub>10</sub> is placed in a reverse micelle, the <sup>51</sup>V chemical shifts show that the V<sub>10</sub> is deprotonated consistent with a decreased proton concentration in the intramicellar water pool. Results indicate that a proton gradient exists inside the reverse micelles, leaving the interior neutral while the interfacial region is acidic.

### I. Introduction

A unique microenvironment for carrying out a variety of chemical and biochemical reactions is found in the polar cores of reverse micelles (RMs). RMs have also been used as model systems for studying various reactions in confinement, for example, biological functions, such as enzymatic reactions<sup>1-5</sup> or micellar catalysis.<sup>6–11</sup> Despite their simplicity, they provide an excellent method to study fundamental effects of confinement. The nature of the pH and ionic concentration in the aqueous core of the micelle assumes particular significance for chemical reactions, such as acid-catalyzed reactions,12,13 electrontransfer reactions,<sup>14–16</sup> and biochemical reactions.<sup>1–5</sup> Thus,

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studies probing the nature of the water pool in RMs interest both chemists and life scientists.

RMs of sodium bis(2-ethylhexyl) sulfosuccinate (AOT) in nonpolar organic solvent can solubilize water into organized structures containing a water pool surrounded by surfactant polar headgroups. Varying the amount of water  $([H_2O]/[AOT] = w_0)$ one can vary the size of water pool.<sup>17</sup> For small RMs,  $w_0 < 10$ , the physical characteristics of the intramicellar water differ substantially from those of bulk water.<sup>18-24</sup> However, as the water pool grows, the properties of the intramicellar water has been reported to approach the properties of bulk water. $^{25-33}$ 

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Researchers have used various molecular probes to investigate properties of the RMs, such as microviscosity,<sup>34,35</sup> pH,<sup>36</sup> and polarity.37 For example, measuring transient absorption of Auramine O, Hirose et al. found substantially slower rotational dynamics inside the RMs than in bulk solution.<sup>34,35</sup> Similarly, magnetic resonance studies of nitroxide probes such as TEMPO or TEMPAMINE show reduced rotational motion when confined in RMs.38 Many different molecular probes have been used to gauge pH in RMs. From steady-state fluorescence and absorption spectroscopy of proton-donating molecules, pyranine and fluorescein, Hasegawa found that AOT headgroups appear to buffer the RM interior.<sup>39</sup> In another study, Biswas et al. showed that pH values at the interface of AOT RMs were greater than those at their cores using absorption and emission spectroscopy of 7-hydroxycoumarin and pyranine.<sup>40</sup> The chemical shift of inorganic phosphate can also be used to indicate the pH of a solution; Fujii et al. have used <sup>31</sup>P NMR spectroscopy to monitor environmental changes in RMs.41,42 These experiments also suggest that the AOT surface buffers the intramicellar aqueous solution. In the work presented here, we utilize a charged inorganic polyoxovanadate, specifically decavanadate, as a probe for the RM water pool.

Vanadium(V) undergoes a range of protonation and oligomerization equilibria in aqueous solution.<sup>43</sup> The simple and colorless vanadate oligomers including vanadate monomer  $(V_1)$ , dimer  $(V_2)$ , tetramer  $(V_4)$ , and pentamer  $(V_5)$  interchange with each other on the millisecond to second time scale in neutral and basic aqueous solution unlike the yellow-orange decavanadate (V<sub>10</sub>), whose interconversion with the other oligomers occurs on a substantially longer time scale.44 Decavanadate, V10  $([V_{10}O_{28}]^{6-})$ , whose structure is shown in Figure 1, is the thermodynamically stable oxovanadate species in aqueous solutions from pH 3 to 6.45 Below pH 2, V10 rapidly hydrolyzes to  $VO_2^+$ ; although not thermodynamically the most stable form above pH 6, V<sub>10</sub> can persist for limited time periods because of its slow hydrolysis into other vanadate oligomers.45 Three different types of vanadium atoms each bound to six oxygen atoms and in a slightly distorted environment comprise each V<sub>10</sub> molecule. Four V<sub>C</sub> atoms lie in an axial position, four V<sub>B</sub> atoms are in the equatorial plane, and two VA atoms lie at the center of the equatorial plane, as shown in Figure 1a.

Vanadium compounds are excellent probe of their surroundings due to the quadrupolar relaxation of vanadium in solution, along with the natural abundance (99.76%) and receptivity of

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*Figure 1.* (a) Structure of  $[V_{10}O_{28}]^{6-}$  (V<sub>10</sub>). The three different types of V atoms are labeled  $V_A$ ,  $V_B$ , and  $V_C$ . (b) Space-filled model of  $V_{10}$ .

the <sup>51</sup>V nucleus from its magnetic moment. <sup>51</sup>V NMR chemical shifts serve as excellent diagnostics to monitor small changes in the vanadium(V) coordination environment as they are very sensitive to changes in the electronic nature of the vanadium atom.46-48 Therefore, <sup>51</sup>V NMR spectroscopy is a convenient method to monitor the speciation of oxovanadates. The distribution of and interaction between oligomers present in the <sup>51</sup>V NMR spectrum can reflect changes in solution pH and ionic strength.43,45,46,49 At pH values from 3 to 6, V10 dominates the spectrum, and its changes in protonation states allow pH changes to be monitored. In its deprotonated form,  $V_{10}$  carries a -6charge; it protonates as the pH drops below 6 as shown in eqs 1 to 3.46

$$V_{10}O_{28}^{6-} + H^+ \rightarrow HV_{10}O_{28}^{5-} \text{ p}K_1 = 5.5 - 6.0^{45}$$
(1)

$$HV_{10}O_{28}^{5-} + H^+ \rightarrow H_2V_{10}O_{28}^{4-} \text{ pK}_2 = 3.1 - 3.7^{45}$$
(2)

$$H_2 V_{10} O_{28}^{4-} + H^+ \rightarrow H_3 V_{10} O_{28}^{3-} \text{ pK}_3 \approx 2^{45}$$
 (3)

At pH 3 the three different vanadium atoms in the trianionic V<sub>10</sub> display <sup>51</sup>V NMR signals at -525, -507, and -424 ppm corresponding to V<sub>C</sub>, V<sub>B</sub>, and V<sub>A</sub>, respectively. Successive deprotonation of the V<sub>10</sub> anion leads to a downfield chemical shifts for the  $V_C$  and  $V_B\ {}^{51}V$  NMR peaks to -516 and -500ppm, respectively. This deprotonation has also been substantiated in detail by <sup>17</sup>O NMR studies.<sup>50</sup> Lifetime and line width measurements further report on the mobility of the probe. Combined, these properties indicate that  $V_{10}$  is an excellent

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*Figure 2.* Representative <sup>51</sup>V NMR spectra of V<sub>10</sub> (a) in stock 10 mM (100 mM V atoms) solution at pH 7.0 and inside reverse micelles created with pH 7.0 stock solution; (b) in 10 mM stock solution at pH 3.1 and inside reverse micelles created with pH 3.1 stock solution. The spectra were recorded at 78.9 MHz using <sup>51</sup>V NMR parameters described previously.<sup>44,75</sup>

spectroscopic probe to explore aqueous solutions, such as the RM interior.<sup>51</sup>

The work presented here utilizes  $V_{10}$  as a probe to explore the nature of the RM interior in large water/AOT/isooctane RMs with  $w_0 = 12$ , 16, and 20. The substantial and uniformly distributed negative charge on the  $V_{10}$  molecule should lead it to reside in the water pool away from the negatively charged headgroups of AOT RMs. Following the <sup>51</sup>V NMR of the intramicellar  $V_{10}$  allows us to measure properties of the micellar interior. While studies suggest that, in large RMs  $w_0 > 10$ , the amount and nature of water in the environment sampled by the  $V_{10}$  should be those of bulk water,<sup>25–33</sup> the results from the experiments reported here indicate otherwise. Results show that the influence of the RM interior leads to deprotonation of  $V_{10}$ introduced in an aqueous solution at pH 3 and that the microviscosity remains higher than that of bulk water.

#### **II. Results**

II. A. <sup>51</sup>V NMR Studies of V<sub>10</sub> in Aqueous Solution and in RMs. Solution preparations are described in the figure captions of data shown, but more detailed experimental information is available in the Supporting Information. We measured the spectrum of V<sub>10</sub> in RM samples generated using stock solutions ranging from pH 3.1 to 8.0. Figure 2 shows the <sup>51</sup>V NMR spectra obtained for  $V_{10}$  in RMs with  $w_0 = 12$ , 16, and 20, and the original stock solution. Three spectral features comprise the  $V_{10}$  spectrum; two peaks appear at -525.3 and -501.0 ppm in the pH 3.1 stock solution differing from those observed at pH 7.0, -514.3 and -498.3 ppm. These chemical shifts reflect the different protonation states of V<sub>10</sub> present at the two pH values. In contrast, the spectra of  $V_{10}$  in the RMs remain surprisingly constant; neither pH nor w<sub>0</sub> appears to alter the chemical shifts. Figure 2 also shows substantial increases in line widths for the three  $V_{10}$  signals in the RMs.



*Figure 3.* <sup>51</sup>V NMR chemical shifts as a function of pH for each vanadium atom type ( $V_C$ ,  $V_B$ , and  $V_A$ ) in  $V_{10}$ . Data were obtained using experimental detail referred to in the caption of Figure 2. Error bars indicate standard error (SE) on triplicate measurements; when none are visible symbols cover error bars.

Given the changes in the chemical shifts observed between pH 3.1 and 7.0 shown in Figure 2, we measured the spectra of  $V_{10}$  in RM samples as a function of stock solution pH from 3 to 8. Figure 3 shows the chemical shifts observed for the three different V atoms present in  $V_{10}$  as a function of the pH. The chemical shifts for  $V_{10}$  in the stock solution are also shown for comparison. The chemical shift of  $V_C$  (see Figure 1) in stock solution at pH 3 differs from the corresponding  $V_C$  in the RM by 8 ppm, while that for the  $V_B$  atom is nearly 6 ppm. In contrast, the chemical shift of  $V_A$  from the stock solution is conserved inside the AOT RM. Most likely because the  $V_A$  atom is deeply buried within  $V_{10}$ , its chemical shift is less sensitive to changes in the external environment. The line widths also change as the stock solution pH is varied (see Figure S1, Supporting Information).

To explore the impact of the RM environment, we present in Figure 4 the chemical shifts of the V atoms in the V<sub>10</sub> molecule as a function of  $w_0$  for RMs formed from V<sub>10</sub> stock solutions at pH 3.1 and 7.0. The chemical shifts for the molecule in the aqueous stock solutions at the two bracketing pH values, 3.1 and 7.0, are shown for comparison. These plots reveal large differences in the chemical shifts for V<sub>C</sub> and V<sub>B</sub> inside the RMs compared to the stock solution at pH = 3.1 for all RM sizes probed. A gradual change in the chemical shift was observed for the entire  $w_0$  range studied. While the chemical shifts observed differ little from one RM size to another when formed with a pH 7.0 stock solution, differences are observed with respect to the stock solution chemical shift for V<sub>B</sub> and V<sub>C</sub>. In

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*Figure 4.* <sup>51</sup>V NMR chemical shifts obtained at 78.9 MHz as a function of  $w_0$  in reverse micelles prepared from AOT (0.5 M)/isooctane and an aqueous 10 mM V<sub>10</sub> stock solution at pH 3.1 (solid squares) and pH 7.0 (open squares) for each vanadium atom type (V<sub>C</sub>, V<sub>B</sub>, and V<sub>A</sub>) in V<sub>10</sub>. <sup>51</sup>V NMR chemical shifts for V<sub>10</sub> in stock solution at pH 3.1 (solid line) and 7.0 (dashed line) are shown for comparison.

RMs containing the pH 3.1 stock solution, the chemical shift is closer to the chemical shift in the stock solution in the larger RMs. However, even at  $w_0 = 20$  the chemical shift is far from that of stock solution, suggesting that these environments are very different.

The line widths of the <sup>51</sup>V signals in the V<sub>10</sub> molecule are shown in Figure 5 as a function of  $w_0$  for RMs formed from pH 3.1 and 7.0 stock solutions. Changes in the line widths for V<sub>B</sub> and V<sub>C</sub> signals are greatest for the RMs prepared from the pH 7.0 stock solution. Specifically, spectral features of V<sub>10</sub> in the stock solution broaden in RMs, with broader signals in smaller RMs. The increasing line width trend is greatest for V<sub>A</sub> in the smallest RMs. Lifetime experiments were done, documenting that the observed changes are attributed to changes in molecular environment. While the line width decreases with increasing RM size, it never reaches the value of the stock solution as can be observed in plots of line width for V<sub>C</sub>, V<sub>B</sub>, and V<sub>A</sub> as a function of  $w_0$  shown in Figure 5. The <sup>51</sup>V NMR line widths also vary as a function of pH but show no clear trends (Figure S1, Supporting Information).

II. B. Dynamic Light-Scattering And Conductivity Experiments Characterizing the RMs Containing Water and  $V_{10}$ . Because addition of  $V_{10}$  to the microemulsion could change the characteristics of the solution, <sup>52–56</sup> we performed a range of experiments to explore the nature of the microemulsions,



*Figure 5.* <sup>51</sup>V NMR line widths obtained at 78.9 MHz as a function of  $w_0$  in reverse micelles prepared from an AOT (0.5 M)/isooctane and a 10 mM V<sub>10</sub> stock solution at pH 3.1 (solid squares) and pH 7.0 (open squares) for each vanadium atom type (V<sub>C</sub>, V<sub>B</sub>, and V<sub>A</sub>) in V<sub>10</sub>. <sup>51</sup>V NMR line widths for V<sub>10</sub> in stock solution at pH 3.1 (solid line) and 7.0 (dashed line) are shown for comparison.

**Table 1.** Reverse Micellar Properties: Radii for Reverse Micelles Containing Aqueous 10 mM V<sub>10</sub> or Pure Water Measured Using Dynamic Light Scattering at a Viscosity of 0.691 cP and a Refractive Index of 1.391 Measured at 826.6 nm and at 25 °C, and Number of V<sub>10</sub> Molecules per RM in the AOT (0.2 M)/ Isooctane System

	000
12 $3.6 \pm 0.1$ $4.0 \pm 0.1$	3
16 $4.0 \pm 0.3$ $4.1 \pm 0.1$	6
20 $4.2 \pm 0.1$ $4.5 \pm 0.2$	11

<sup>*a*</sup> Stock solution  $pH = 6.0^{-b}$  Stock solution  $pH = 5.7^{-c}$  Estimated from the reverse micelle aggregation number<sup>57</sup> and overall vanadate concentration.

including dynamic light scattering, to compare the sizes of RMs containing water and  $V_{10}$  solution, as well as conductivity and viscosity measurements. The sizes of RMs could be measured with an overall AOT concentration of 0.2 M where the viscosity of the solution is close to that for isooctane. Data for these measurements are given in Table 1. Also given are the estimated numbers of  $V_{10}$  probe molecules present in the RMs. These data show that the size of the RMs varies within the error of the measurement. Thus, within the parameters of the experiments

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performed here, our results show that the RMs are not perturbed by dissolution of the  $V_{10}$  molecule.

Electrical conductivity of the RM system was also used to compare RMs containing water and V<sub>10</sub> solutions. The conductivity of water-in-oil microemulsions in RMs containing water shows variations over many orders of magnitude as the phase changes.58,59 The conductivity of RMs encapsulating V10 prepared from 0.5 M AOT stock solution was 0.4  $\mu$ S cm<sup>-1</sup>. Corresponding RMs prepared with 0.5 M AOT stock solution and pure water show the same conductivity. Similarly, the electrical conductivity of RM solutions prepared from 0.2 M AOT is 0.1  $\mu$ S cm<sup>-1</sup> whether pure water or aqueous V<sub>10</sub> forms the interior. In contrast, the electrical conductivity of the 100 mM NaVO<sub>3</sub> aqueous stock solution is significantly higher, 8.7 mS cm<sup>-1</sup>.

The studies described above were all obtained for RMs formed in 0.5 M AOT solutions that show somewhat increased conductivity compared to lower AOT concentrations. Our previous work suggests that at 0.5 M AOT in isooctane, RMs tend to aggregate or flocculate.<sup>51</sup> Thus, we explored the impact of AOT concentration to learn what effect, if any, the RM concentration has on the observed <sup>51</sup>V NMR signals. Upon dilution of the AOT to 0.2 M, similar <sup>51</sup>V NMR experiments yielded results indistinguishable from those carried out at the higher AOT concentration. The only differences observed were the change of vanadate oligomer speciation toward the neutral pH and the poorer signal-to-noise resulting from lower overall concentration of V10. Therefore, although the RMs may aggregate or flocculate in 0.5 M solutions, the environment of the V<sub>10</sub> species is conserved. These results indicate that incorporation of V<sub>10</sub> in the RMs does not substantially perturb the RM structure.

#### **III. Discussion**

In AOT RMs, water molecules near the interface can solvate AOT headgroups, Na<sup>+</sup> counterions, or they can interact with other water molecules. Thus, it is not surprising that results from a wide range of experiments show that water near interfaces behaves differently than it does in bulk solution.<sup>17,23,24,33,60-66</sup> Water molecules perturbed by interactions with AOT and bulklike water equilibrate continuously as molecules move from interface to core and back. However, in large RMs,  $w_0 > 10$  or more generally when the water pool radius exceeds 22 Å, the intramicellar water is reported to develop characteristics similar to those of bulk water that can dominate the results observed.<sup>25-33</sup> Thus, in experiments utilizing large ( $w_0 \ge 10$ ) RMs, the water often displays properties similar or identical to those observed

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in bulk solution. When results diverge from bulk water characteristics, researchers generally suggest that the molecular probes utilized must reside near the RM interface.<sup>67</sup> The results reported in this paper show that, even in the "bulk-like" water pools found in large RMs, RM water properties can differ substantially from those of bulk water.

All experiments reported here have been performed on RMs with  $w_0 > 10$ , when the RMs are large enough to support a bulk-like water pool. The large charge of the  $V_{10}$  molecule and its V-O bonds make this molecule very hydrophilic. That, combined with Coulombic repulsion between the substantial negative charge on  $V_{10}$  (between -6 and -4) and the negatively charged AOT headgroups, should drive the V<sub>10</sub> into the RM water pool. The possibility still exists that  $V_{10}$  interferes with the formation and nature of the RMs, and the presence of the V<sub>10</sub> molecule in the RM interior could itself influence the environment. On the basis of its size, we estimate that one  $V_{10}$ molecule takes up the same amount of volume as approximately 11-12 water molecules. Given that there are thousands of water molecules inside these RMs, it seems unlikely that removing 11 water molecules per  $V_{10}$  should strongly perturb the system. We explored the impact of V10 on RM formation and character using dynamic light-scattering experiments and conductivity. Dynamic light scattering experiments showed that the hydrodynamic radii of RMs and polydispersity for RMs with  $w_0 =$ 12, 16, and 20 formed with water or aqueous stock solutions of  $V_{10}$  are the same (Table 1). This indicates that solubilization of the  $V_{10}$  molecule in the RMs changes neither their size nor their shape. Furthermore, RMs formed with pure water or aqueous V<sub>10</sub> stock solution over the range of pH values show similar conductivity, 0.4  $\mu$ S cm<sup>-1</sup> and 0.1  $\mu$ S cm<sup>-1</sup>, at 0.5 and 0.2 M AOT, respectively. At concentrations above 0.2 M AOT the RMs flocculate, but <sup>51</sup>V NMR experiments suggest that the environment of V10 remains similar for 0.5 and 0.2 M AOT concentration. Electrical conductivity measurements show that addition of  $V_{10}$  in the water pool does not change the microemulsion phase. These results leave us confident that the V10 molecule does not interfere with RM formation or character and that it does not insert into the RM interface as we have observed for other vanadium-containing probes.51

The location of the V<sub>10</sub> molecule in the RM should influence the observed <sup>51</sup>V NMR signal. Indeed, there is substantial evidence that the environment near the micellar interface differs significantly from the micellar core.<sup>21,23,24,68</sup> Sampling the interfacial region should lead the V10 molecule to report on water with properties different from bulk. At higher water content,  $w_0 > 10$ , the interfacial properties in AOT RMs formed in aliphatic hydrocarbon solvents, such as isooctane in the studies discussed here, become independent of water content.<sup>17,68,69</sup> Molecular probes sensing microviscosity or micropolarity that reside at the RM interface show significant variations with increasing RM size but approach a constant value for larger RMs. While Coulombic repulsion is not always sufficient to determine molecule location in a RM,51,70 the continuing changes

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Figure 6. (a) Geometric depiction of the location of the  $V_{10}$  complex in a reverse micelle corresponding to parameters from eqs 4-8. (b) Free energy as a function of intercenter distance, s, as described in the text. Parameters for graph are as follows,  $R_{V10} = 7$ ,  $R_{RMs} = 75$ ,  $V_{10}$  charge = -6, and RMs charge is set to -1.

in chemical shifts and line widths (Figures 4 and 5) that we observe suggest  $V_{10}$  does not remain solvated at the RM inner interface and rather resides in the core water pool. In addition, we note that in small AOT RMs, dramatic changes to <sup>51</sup>V NMR signals are apparently consistent with the  $V_{10}$  molecule approaching and interacting with the interface. (We have measured <sup>51</sup>V NMR spectra for V<sub>10</sub> in AOT RMs with  $w_0 < 10$ . In these systems, we observe unequal influence on the peaks in the NMR spectra as V<sub>B</sub> and V<sub>C</sub> atoms interact more strongly with the interface. These data will be published in a forthcoming paper.)

Because the  $V_{10}$  molecule carries a significant negative charge and the RM interface is also charged, the simple model system of two charged, eccentric spheres developed by Sengupta and Papadopoulos<sup>71</sup> can be used to model the  $V_{10}$  in the RMs. This model places a rigid, solid, inner sphere, that is, the  $V_{10}$ molecule, inside a hollow, rigid, outer sphere, the RM, as depicted in Figure 6a. No shape fluctuations are permitted for either sphere. While the form of the  $V_{10}$  molecule is not perfectly spherical, as shown in Figure 1b, a spherical form with uniform charge density is a reasonable first approximation. The  $V_{10}$ (inner sphere) is allowed to move radially from the RM (outer sphere) center to the interface; the displacement of RM and  $V_{10}$ centers is given by distance s. Then using the Poisson-Boltzmann equation to determine the electric potential,  $\Psi$ , eq 4,

$$\nabla^2 \Psi = \frac{8\pi nez}{\epsilon} \sinh \frac{ze\Psi}{kT} \tag{4}$$

where n is the number concentration of ions in the bulk, e is the charge of the electron, z is the valency of ions,  $\epsilon$  is the dielectric constant of the electrolyte, k is the Boltzmann constant, and T is the temperature, we calculate the free energy of interaction between the  $V_{10}$  and the RM,<sup>72</sup>

$$\Delta F = F_{\rm sys} - \sum_{i=1}^{2} F_i \tag{5}$$

where

$$F_{\text{tot}} = -\frac{\epsilon}{8\pi} \int \int \int_{\nu} |\nabla^2 \Psi|^2 \, \mathrm{d}V - 2nkT \int \int \int_{\nu} \left(\cosh\frac{ze\Psi}{kT} - 1\right) \mathrm{d}V$$
(6)

 $F_{\text{tot}}$  can be simplified and made dimensionless by the substitution

$$F_{\rm tot}^* = \frac{F_{\rm tot}\kappa^3}{nkT} \tag{7}$$

where

$$\kappa^2 = \frac{8\pi n e^2 z^2}{\epsilon kT} \tag{8}$$

Here  $F_{sys}$  is the energy when the surfaces of the spheres interact, and  $F_i$  is the energy when the spheres have no interactions. To compute  $\Delta F$  requires the free energies of the individual spheres that depend on a few parameters, such as surface charge and sphere radii that we estimate from data about AOT RMs and the V<sub>10</sub>.<sup>73</sup> This allows  $\Delta F^*$  to be determined as a function of the separation distance of the sphere centers, s, defined in Figure 6a. (We made a few simplifying assumptions to compute  $\Delta F^*$ . Applying the Poisson-Boltzmann equation assumes a symmetric electrolyte solution, which does not model the RM interior perfectly. However, the calculation is valid as long as the density of ions at the interfaces remains fixed.<sup>74</sup> In addition, as the dielectric constant of the particle is smaller than that of the electrolyte solution, the overall result is a constant Debye length,  $\kappa$ , leading to overall scaling by  $\kappa$ .<sup>75</sup>) While a uniform charge density is assumed for each sphere, when we varied the charges on the inner and outer spheres over a wide range of values, we obtained comparable results.

Figure 6b shows  $\Delta F^*$ , the dimensionless interaction energy, as a function of the separation distance. At s = 0, the V<sub>10</sub> resides at the center of the RM; s = 12 places V<sub>10</sub> at the interface. For plots shown in Figure 6b, we assume that the  $V_{10}$  radius is 9.3%

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of the RM radius, on the basis of the V<sub>10</sub> crystal structure<sup>73</sup> and a RM with  $w_0 = 12$ . Varying the RM size from 12 to 20 led to changes too small to observe. Variation of the interfacial charge between -1 and -2, corresponding to the anticipated number of AOT headgroups with which the V<sub>10</sub> complex interacts, leads to very minor differences in  $\Delta F$  or  $\Delta F^*$ .

 $V_{10}$  sequestered in the AOT RM allows us to compare different parameters and how they affect the interaction of  $V_{10}$ at the interface. The plot shown in Figure 6b indicates that the free energy increases with increasing *s*, which suggests that if Coulombic interactions dominate the placement, the  $V_{10}$  should reside preferentially away from the micellar interface. Only if the  $V_{10}$  possesses an asymmetric charge distribution or if its size approaches the size of the interfacial cavity can it reside near or in the interfacial layer. On the basis of data about  $V_{10}^{73}$ and the protonation state we measure with NMR, the  $V_{10}$ complex should possess a symmetric charge distribution therefore placing it in the RM water pool and not at the interface.

Both literature precedent<sup>69,71</sup> and our calculations indicate that the large, negatively charged  $V_{10}$  ( $[V_{10}O_{28}]^{6-}$ ) molecule should reside away from the interface where it would probe a bulklike water environment. Preference for this location should be greater the greater the charge of the  $V_{10}$  molecule. Varying the pH of the stock solution from 3 to 8 results in a net change from -6 to -4 on the  $V_{10}$  molecule. Thus, one might expect these experiments to yield spectra similar or identical to those measured in bulk aqueous solution, but the data indicate otherwise.

The  $V_{10}$  signals both shift downfield, and the line widths narrow as the water content  $w_0$  increases. However, neither chemical shift nor line width in any of the systems at any starting pH value achieves the corresponding values we observe for the molecule in bulk solution (Figures 4 and 5). Many studies reported in the literature indicate that a bulk-like water pool should have developed at  $w_0 > 10.25-32$  The continuing changes we observe as  $w_0$  increases from 12 to 20 and the fact that the bulk values are never reached show that the  $V_{10}$  environment in all the RMs studied differs from that of bulk water and continues to change in these systems. Thus, we believe that it is the water that has not reached its bulk limit rather than a strong perturbation of the vanadium stock solution on the water. The fact that the chemical shifts and the line widths continue to change above this hydration level suggests that the free water pool has still not achieved bulk-like character with respect to the  $V_{10}$  environment. Thus, if the intramicellar water pool has the properties of bulk water, then the V<sub>10</sub> NMR signals should reflect these water pool properties and should display, at most, minor differences from those observed in bulk solution.

The chemical shifts of the surface-exposed sites  $V_C$  and  $V_B$  in RMs differ dramatically from those observed for  $V_{10}$  in aqueous solution, especially for low pH stock solutions. The chemical shifts of protonated  $V_{10}$  in aqueous stock solution at pH 3.1 are -525.3 ( $V_C$ ), -507.4 ( $V_B$ ), and -425.3 ( $V_A$ ), whereas the chemical shifts of this solution added to the RMs appear at -517.0 ( $V_C$ ), -501.6 ( $V_B$ ), and -424.1 ( $V_A$ ). Given its sensitivity to protonation,<sup>47,48</sup> the <sup>51</sup>V chemical shift of  $V_{10}$  can be used to probe the local pH of environment inside the RMs. The chemical shift differences for  $V_C$  and  $V_B$  are consistent with deprotonation of  $V_{10}$  in bulk solution.<sup>48</sup> This observation suggests that proton concentration in the environ-

ment sensed by  $V_{10}$  in the RMs is lower than that of the stock solution from which the RMs were formed. Interestingly, the chemical shifts we observe for  $V_{10}$  inside the RMs appear conserved regardless of the protonation state in the stock solution (Figures 2 and 3, Table S1, Supporting Information).

The smallest changes in chemical shift and line width are observed for the  $V_A$  atom in the  $V_{10}$  molecule (Figures 3 and 5). This nonoxovanadium atom differs significantly from the surface-exposed  $V_B$  and  $V_C$  atoms. The chemical shift for  $V_A$  is similar in RMs and in bulk solution at all pH values. Completely surrounded by oxygen and other vanadium atoms, its environment isolates  $V_A$  from the exterior, rendering it far less sensitive to changes in the environment. Thus, this vanadium serves as an excellent internal reference for this probe. In other words, changes observed for  $V_B$  and  $V_C$  but not for the  $V_A$  atom reflect surface-type interactions of  $V_B$  and  $V_C$ .

The conundrum arises, why does the  $V_{10}$  molecule report different proton concentration in the RM interior than the starting stock solution? The significant charge on the  $V_{10}$  molecule should lead it to reside well solvated by many intramicellar water molecules while sodium counterions balance both the charge of the  $V_{10}$  and the AOT sulfonate headgroups. However, the <sup>51</sup>V NMR chemical shifts indicate that the RM interior differs substantially from the original stock solution. An inhomogeneous electrical potential can arise inside the RM when counterions dissociate from the interfacial region.<sup>76</sup> A theoretical calculation shows a significant decrease in the potential, increasing the dielectric constant of the media as  $w_0$  increases which could determine the distribution of charged species inside the RMs.<sup>76,77</sup> Recent calculations modeling ion exchange inside the RM show that if the diameter of an added cation is greater than the Na<sup>+</sup> counterion diameter, the cation migrates toward the headgroup region of the RM.76 As protons aggregate to H<sub>3</sub>O<sup>+</sup> and larger aggregates in water,<sup>78-80</sup> their size exceeds the Na<sup>+</sup> and should lead H<sup>+</sup> to migrate preferentially to the interface while Na<sup>+</sup> migrates to the micellar interior as depicted in Figure 7. This migration leads to a proton gradient inside the RM giving rise to different apparent pH values in the interface and in the central water pool. Thus, a complex like  $V_{10}$  located in the water pool of the RM senses a proton concentration that is substantially different than that of the water in the interface or in the original stock solution. Specifically, a symmetrical probe located in the water pool will report a local pH more basic than the original acidic stock solution used for preparation of the RMs. As wo increases, the dielectric potential drops, resulting in a proton migration to the interface, rendering the water pool slightly more acidic than at lower  $w_0$ .

The observation that  $V_{10}$  yields the same <sup>51</sup>V NMR chemical shifts when added to the RMs regardless of the pH of the initial stock solution can be interpreted as the RMs having a buffering effect on the water pool. Our observation that the chemical shifts change toward the values in the bulk stock solution with increasing  $w_0$  increases at a fixed pH supports this interpretation.

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Figure 7. Schematic depiction of the location of  $V_{10}$  with respect to the reverse micellar interface including counterions.

The possibility that the water pool of the RMs is buffered has previously been suggested on the basis of studies carried out with emission spectroscopy.<sup>35,39</sup> The buffering capacity of the water pool in n-heptane/AOT/water RMs was related to the considerably high AOT sulfonate concentration localized at the interface of the aggregate.

The chemical shifts observed for the  $V_{B}$  and  $V_{C}$  atoms in V10 show almost no dependence on the pH of the stock solution used to create the RMs, Figure 3. However, the line width changes as a function of pH (Table S2 and Figure S1, Supporting Information) consistent with changes in microfluidity in the RM interior. Microviscosity has been found to vary with the value of  $w_0$ .<sup>34,81</sup> In small RMs and low  $w_0$ , the water molecules hydrate the polar headgroup of the interface and the accompanying counterions. The key hydrogen bond structure present in the bulk water is destroyed when water interacts with the polar head of the surfactant. Only above  $w_0 = 10$  when the interface has obtained the water molecules needed to complete the hydration sphere will the water molecules begin to form a water pool with characteristics of the "free" bulk-phase water, as determined by properties such as microviscosity and micropolarity.<sup>37,40,68,82</sup> The bound water found at interfaces at very small  $w_0$ , shows a very high microviscosity, and low polarity yields a very rigid interface. As the size of the RM increases the water content and  $w_0$  increases, the microviscosity and, hence, the rigidity of the RM decrease.<sup>37,40,68,82-84</sup> Thus, the water pool and the interface of the RM are decreased in viscosity and increased in mobility as the  $w_0$  increases. Since the line widths for two out of three V signals decrease with increasing  $w_0$ , the complex must sense a less viscous environment as the water content increases. However, it never reaches the value found in bulk water. This result is contrary to the many reports that characterize the water in RMs with  $w_0 > 10$  as bulk water,<sup>52</sup> and suggests that, although by many criteria one would characterize the water as similar to bulk water, these studies show that other properties exist for which the water inside the RM is different than bulk water.

## **IV. Conclusions**

The highly charged inorganic decavanadate anion  $(V_{10})$  is used to probe the water environment in RMs formed by AOT in isooctane. The versatile inorganic probe utilized here, that is V<sub>10</sub>, provides several spectroscopic handles that allow us to probe more than one property of AOT RMs at a time. Due to the large negative charge on V<sub>10</sub> and the negatively charged surfactant headgroups, it is generally presumed that  $V_{10}$  would reside in the RM water pool. Continuing but small changes in the <sup>51</sup>V NMR chemical shifts and line widths with growing RM size confirm this presumption.

The  $V_{10}$  molecule is an effective and unique probe that allows us to measure two very important properties of the intramicellar water pool, that is proton concentration, or local pH, and microviscosity. The results presented here show that, despite sequestering large amounts of water within, the water in these RMs never reaches the values found for bulk water. Furthermore, the  $V_{10}$  molecule shows that the core region of the RMs remains at an apparent pH near neutral. Protons migrate toward the RM interfacial region, leaving the core region with counterions but no excess protons. These results could have tremendous impact for chemical reactions occurring in RMs, such as acid-catalyzed reactions, or biochemical reactions, as enzymatic reactions. While it is not clear that the nature of confined water pools share similarities with the AOT RMs, if they do, the implications could be far reaching for a wide range of processes occurring in confined environments.

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Supporting Information Available: Experimental details. This material is available free of charge via the Internet at http://pubs.acs.org.

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